

Cosmology and the origin of structure

<http://home.fnal.gov/~rocky/>

Rocky I: The universe observed

Rocky II: The growth of cosmological structures

Rocky III: Inflation and the origin of perturbations-1

Rocky IV: Inflation and the origin of perturbations-2

Rocky V: Dark matter

Xth Brazilian School of Cosmology and Gravitation

August 2002, Mangaratiba

Rocky Kolb

Fermilab, University of Chicago, & CERN

Rocky II: Growth of structure

- Linear regime: quantitative analysis

Jeans analysis

Sub-Hubble-radius perturbations (Newtonian)

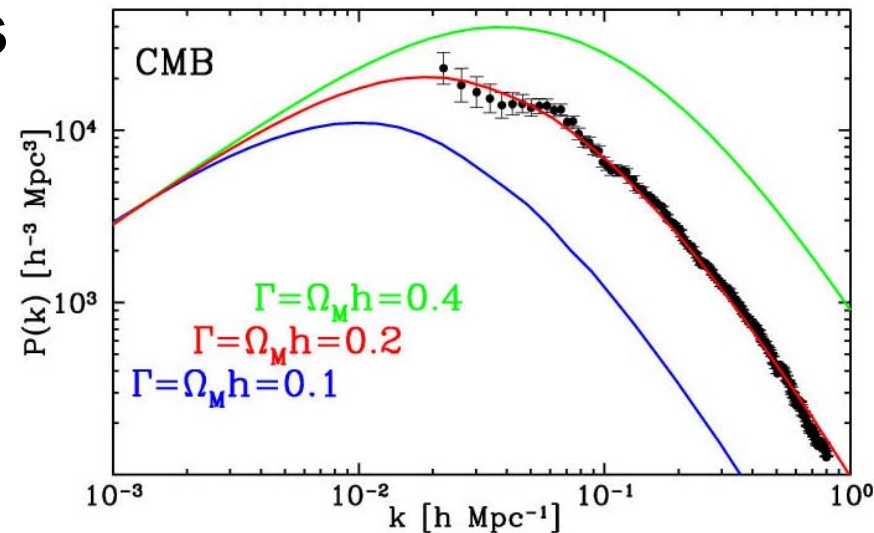
Super-Hubble-radius perturbations (GR)

Harrison-Zel'dovich spectrum

Dissipative processes

The transfer function

Linear evolution



- Non-linear regime: word calculus

Comparison to observations

A few clouds on the horizon

Growth of small perturbations

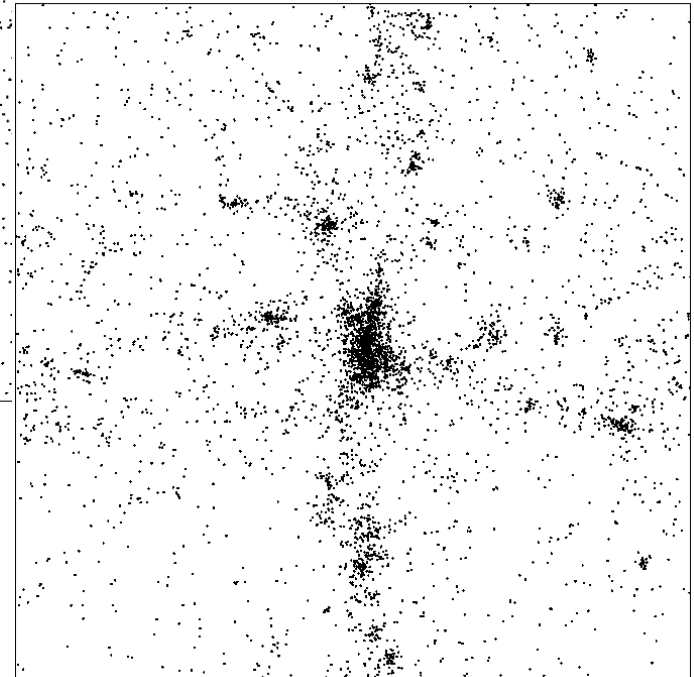
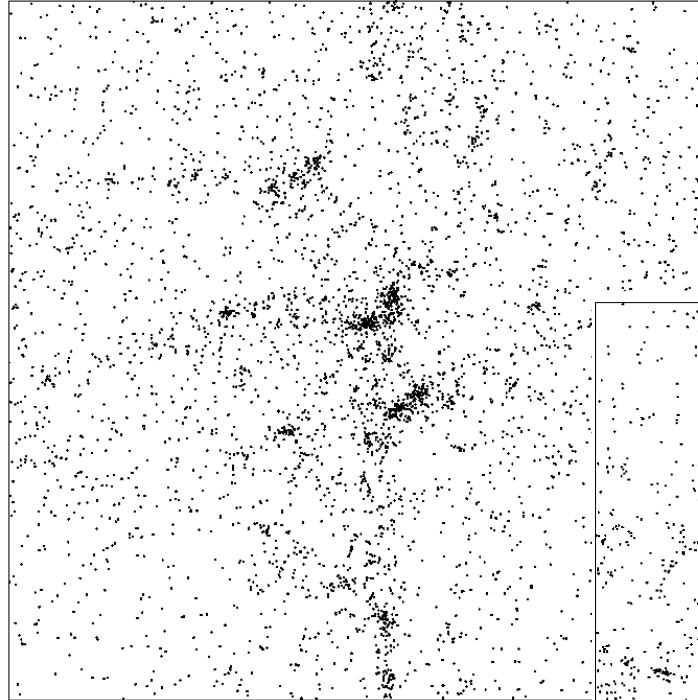
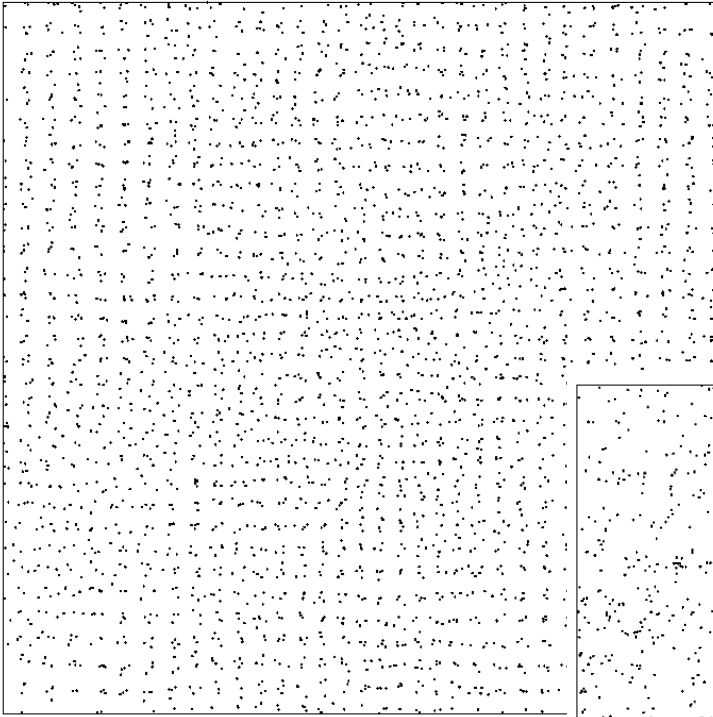
Today (12 Gyr AB)

- radiation and matter decoupled
- $\Delta T/T \sim 10^{-5}$
- $\Delta\rho_G/\rho_G \sim 10^{+6}$

Before recombination (300 kyr AB)

- radiation and matter coupled
- $\Delta T/T \sim 10^{-5}$
- $\Delta\rho_G/\rho_G \sim 10^{-5}$

Seeds of structure



time →

Simulation

Simulation (simiˈleɪʃən). ME. [a. OF., ad. L. *simulationem*.] 1. The action or practice of simulating, with intent to deceive; false pretence, deceitful profession ME. b. Unconscious imitation 1870. 2. A false assumption or display, a surface resemblance or imitation, *of* something.

Oxford English Dictionary

Power spectrum

- Assume there is an average density $\bar{\rho}$
- Expand density contrast $\delta(\vec{x})$ in Fourier modes

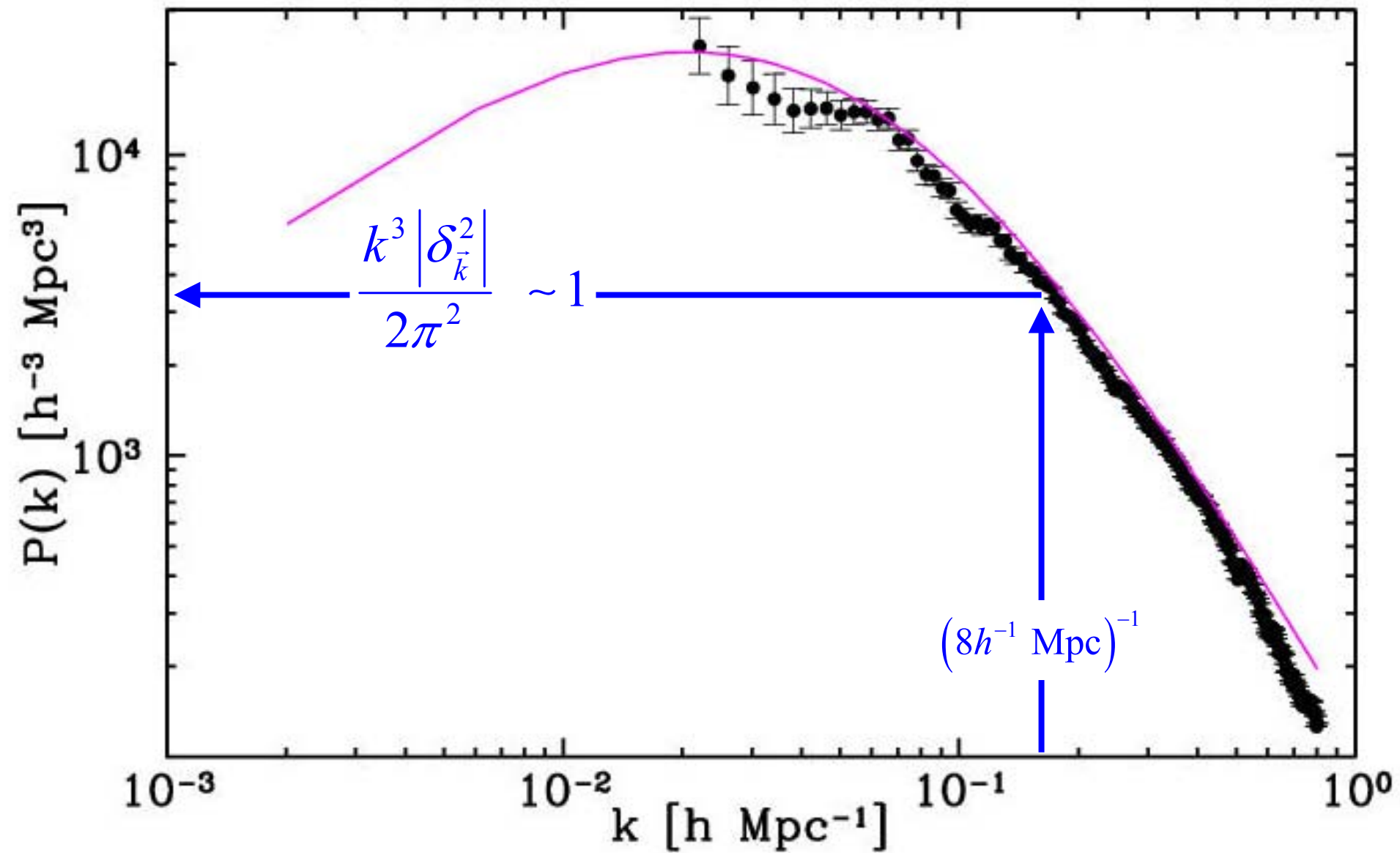
$$\delta(\vec{x}) \equiv \frac{\rho(\vec{x}) - \bar{\rho}}{\bar{\rho}} = \int \delta_{\vec{k}} \exp(-i\vec{k} \cdot \vec{x}) d^3k$$

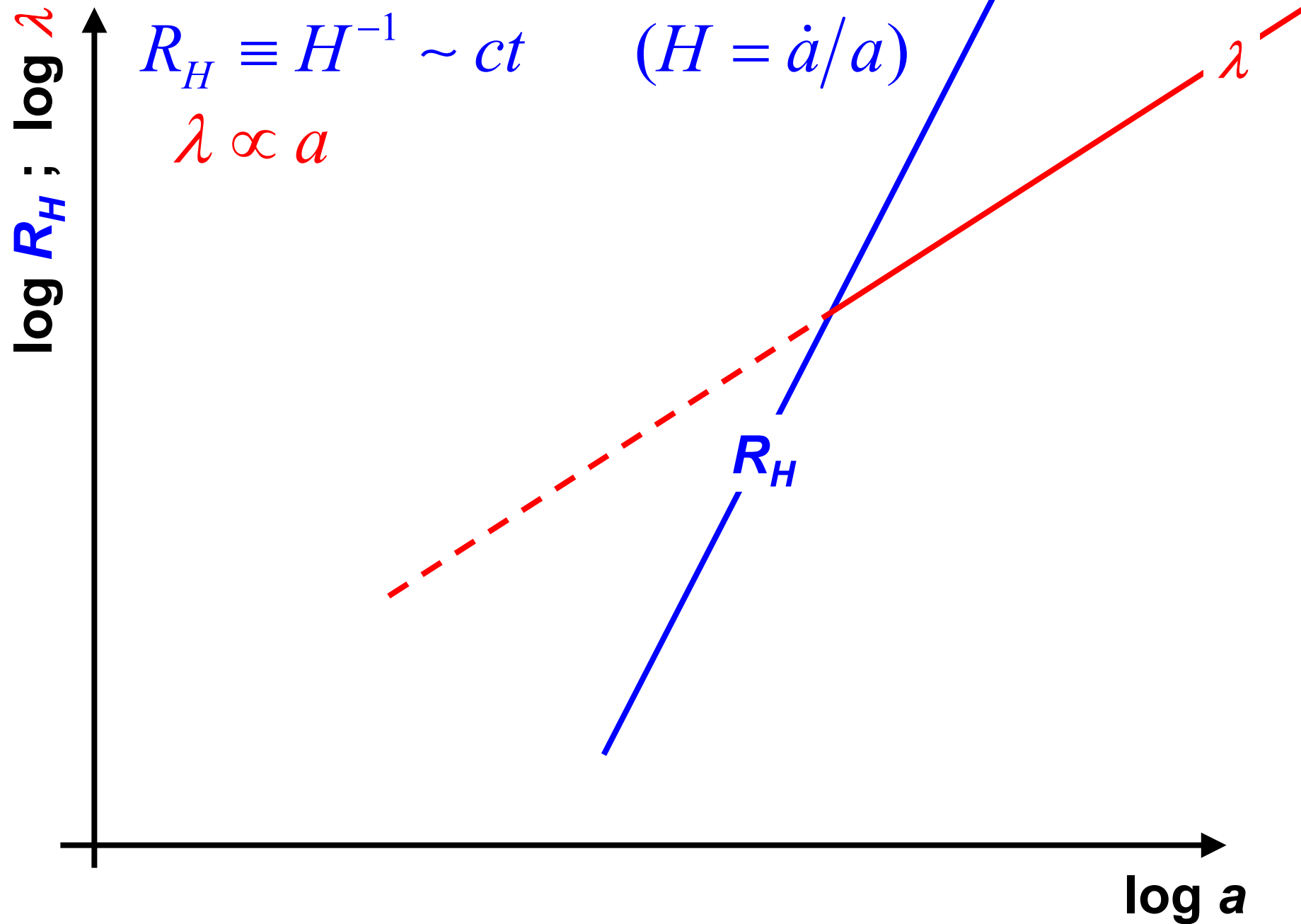
- Autocorrelation function defines power spectrum

$$\left\langle \frac{\delta\rho(\vec{x})}{\rho} \right\rangle^2 = \langle \delta(\vec{x})\delta(\vec{x}) \rangle = \int_0^\infty \frac{dk}{k} \frac{k^3 |\delta_{\vec{k}}^2|}{2\pi^2}$$

$$\Delta^2(k) \equiv \frac{k^3 |\delta_{\vec{k}}^2|}{2\pi^2} \quad P(k) \equiv |\delta_{\vec{k}}^2|$$

Power spectrum







Jeans analysis

Jeans analysis in a non-expanding fluid:

matter density

ρ

pressure

p

velocity field

\vec{v}

gravitational potential

ϕ

$$\left\{ \begin{array}{l} \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0 \\ \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} + \frac{1}{\rho} \vec{\nabla} p + \vec{\nabla} \phi = 0 \\ \nabla^2 \phi = 4\pi G \rho \end{array} \right.$$

Perturb about solution*

$$\rho = \rho_0 = \text{constant}$$

$$\rho = \rho_0 + \rho_1$$

$$p = p_0 = \text{constant}$$

$$p = p_0 + p_1$$

$$\vec{v} = 0 = \text{constant}$$

$$\vec{v} = \vec{v}_0 + \vec{v}_1$$

$$\phi = \phi_0 = \text{constant}$$

$$\phi = \phi_0 + \phi_1$$

$$\frac{\partial^2 \rho_1}{\partial t^2} - v_s^2 \nabla^2 \rho_1 = 4\pi G \rho_1$$

$$v_s^2 = p_1 / \rho_1$$

Jeans analysis

$$\frac{\partial^2 \rho_1}{\partial t^2} - v_s^2 \nabla^2 \rho_1 = 4\pi G \rho_1$$

solutions of $\rho_1(\vec{r}, t) = \delta(\vec{r}, t) \rho_0$
the form

$$= A_k \exp(-i\vec{k} \cdot \vec{r} + i\omega t)$$

ω and k satisfy the dispersion relation $\omega^2 = v_s^2 k^2 - 4\pi G \rho_0$

ω real: perturbations oscillate as sound waves

ω imaginary: exponentially growing (or decaying) modes

Jeans wavenumber $k_J = \left(\frac{4\pi G \rho_0}{v_s^2} \right)^{1/2}$ $k > k_J$ perturbation oscillates
 $k < k_J$ perturbation grows

Jeans mass $M_J = \frac{4\pi}{3} \left(\frac{\pi}{k_J} \right)^3 \rho_0$ $M < M_J$ perturbation oscillates
 $M > M_J$ perturbation grows

gravitational pressure vs. thermal pressure



Sub-Hubble-radius ($R_H = H^{-1}$)

Jeans analysis in an expanding fluid:
scale factor $a(t)$ describes expansion,
unperturbed solution:

$$\rho_0 = \rho_0(t_0) a^{-3}(t) \quad \vec{v}_0 = \frac{\dot{a}}{a} \vec{r} \quad \vec{\nabla} \phi_0 = \frac{4\pi G \rho_0}{3} \vec{r}$$

$$\ddot{\delta}_k + 2 \frac{\dot{a}}{a} \dot{\delta}_k + \left(\frac{v_s^2 k^2}{a^2} - 4\pi G \rho_0 \right) \delta_k = 0$$

- Solution is some sort of Bessel function:
growth or oscillation depends on Jeans criterion
- In matter-dominated era $\rho_0 = (6\pi G t^2)^{-1}$ and $\dot{a}/a = 2/3t$
- For wavenumbers less than Jeans

$$\delta_+(t) = \delta_+(t_i) (t/t_i)^{2/3} \quad \delta_-(t) = \delta_-(t_i) (t/t_i)^{-1}$$



Super-Hubble-radius ($R_H = H^{-1}$)

$$g_{\mu\nu}(\vec{x}, t) = g_{\mu\nu}^{FRW}(t) + \delta g_{\mu\nu}(\vec{x}, t)$$

$$T_{\mu\nu}(\vec{x}, t) = T_{\mu\nu}^{FRW}(t) + \delta T_{\mu\nu}(\vec{x}, t)$$

$$\delta R_{\mu\nu} - (1/2)\delta[g_{\mu\nu}R] = 8\pi G\delta T_{\mu\nu}$$

- complete analysis not for the faint of heart
- interested in “scalar” perturbations
- fourth-order differential equation
- only two solutions “physical”
- other two solutions are “gauge modes”
which can be removed by a coordinate
transformation on the unperturbed metric

$$\delta R_{\mu\nu} - (1/2) \delta [g_{\mu\nu} R] = 8\pi G \delta T_{\mu\nu}$$

Bardeen 1980

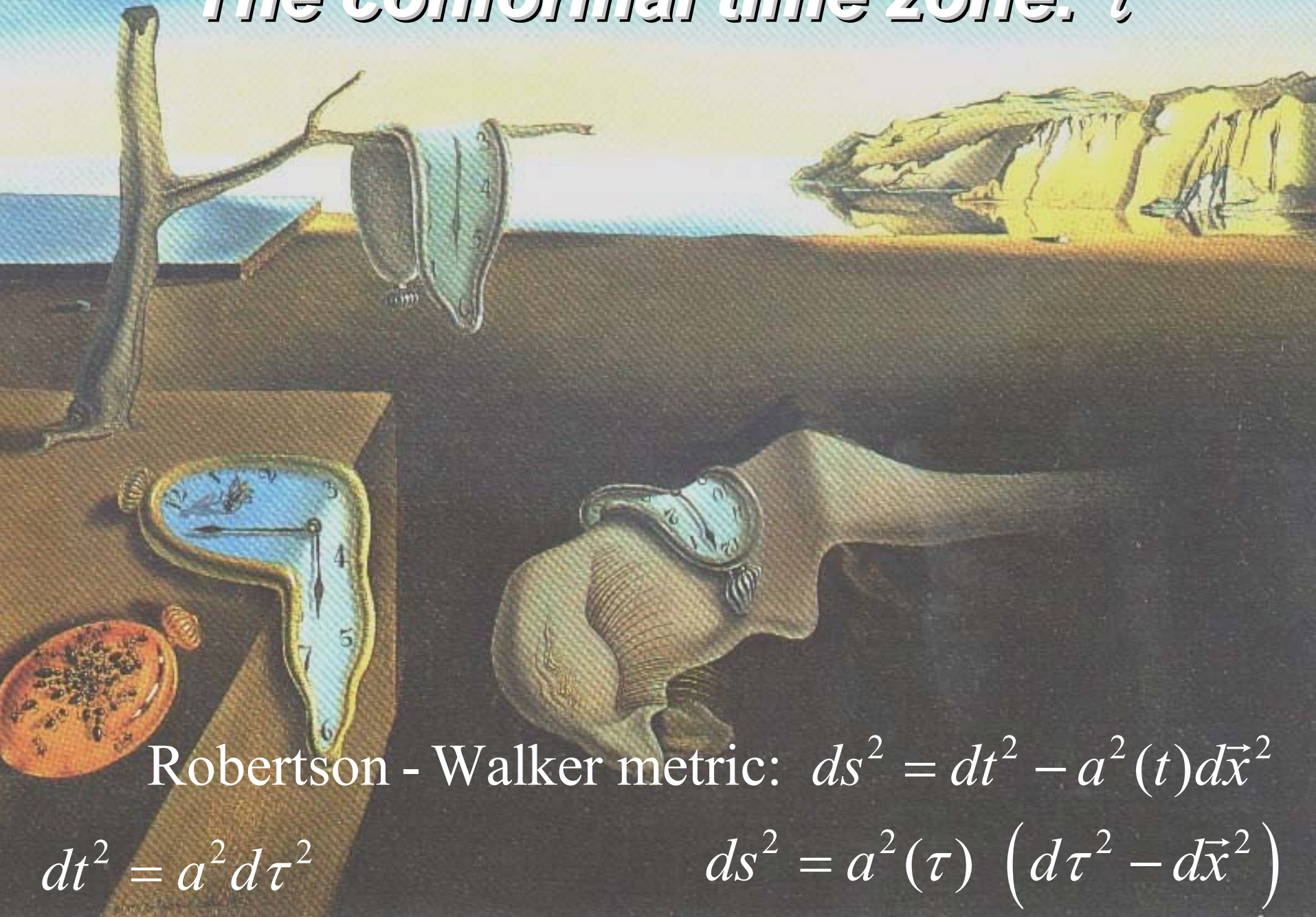
Reference spacetime: flat FRW

$$ds^2 = a^2(\tau) \{ d\tau^2 - \delta_{ij} dx^i dx^j \}$$

$\tau =$ conformal time

$$dt^2 = a^2(\tau) d\tau^2$$

The conformal time zone: τ



Robertson - Walker metric: $ds^2 = dt^2 - a^2(t)d\vec{x}^2$

$$dt^2 = a^2 d\tau^2$$

$$ds^2 = a^2(\tau) \left(d\tau^2 - d\vec{x}^2 \right)$$

$$\delta R_{\mu\nu} - (1/2) \delta [g_{\mu\nu} R] = 8\pi G \delta T_{\mu\nu}$$

Bardeen 1980

Reference spacetime: flat FRW

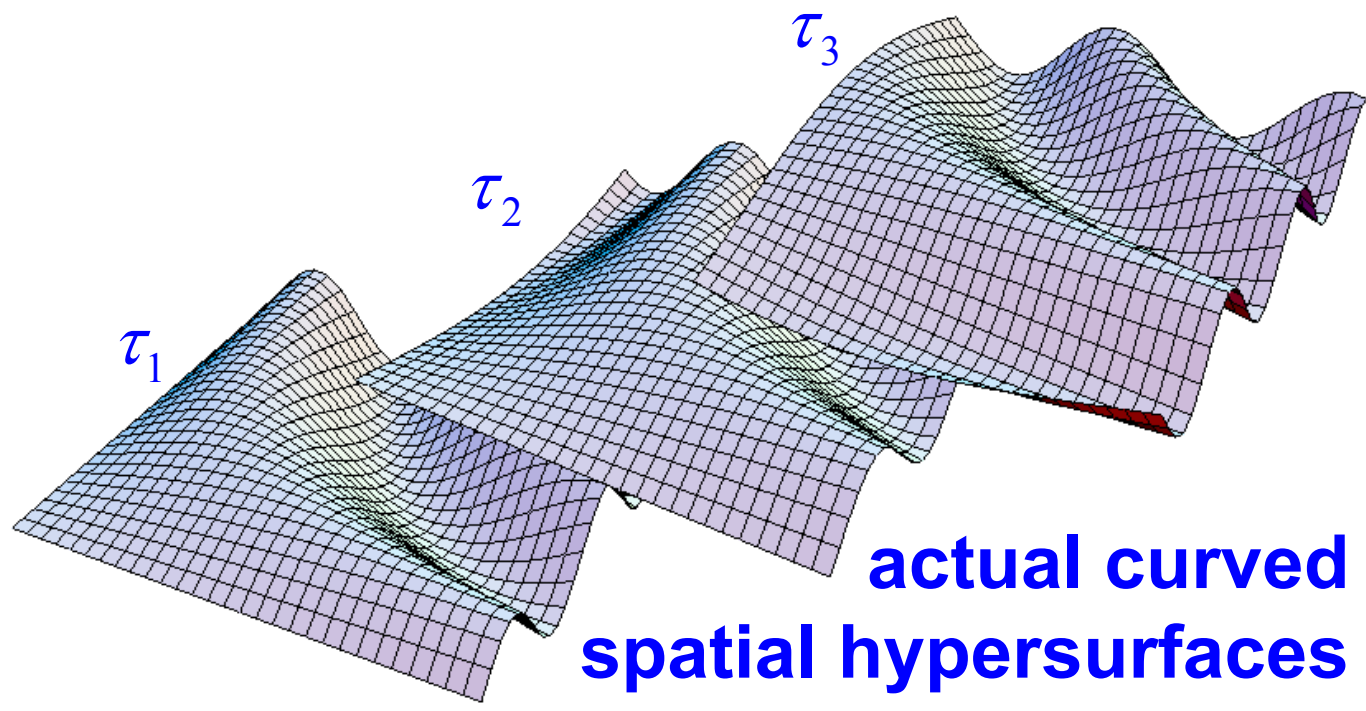
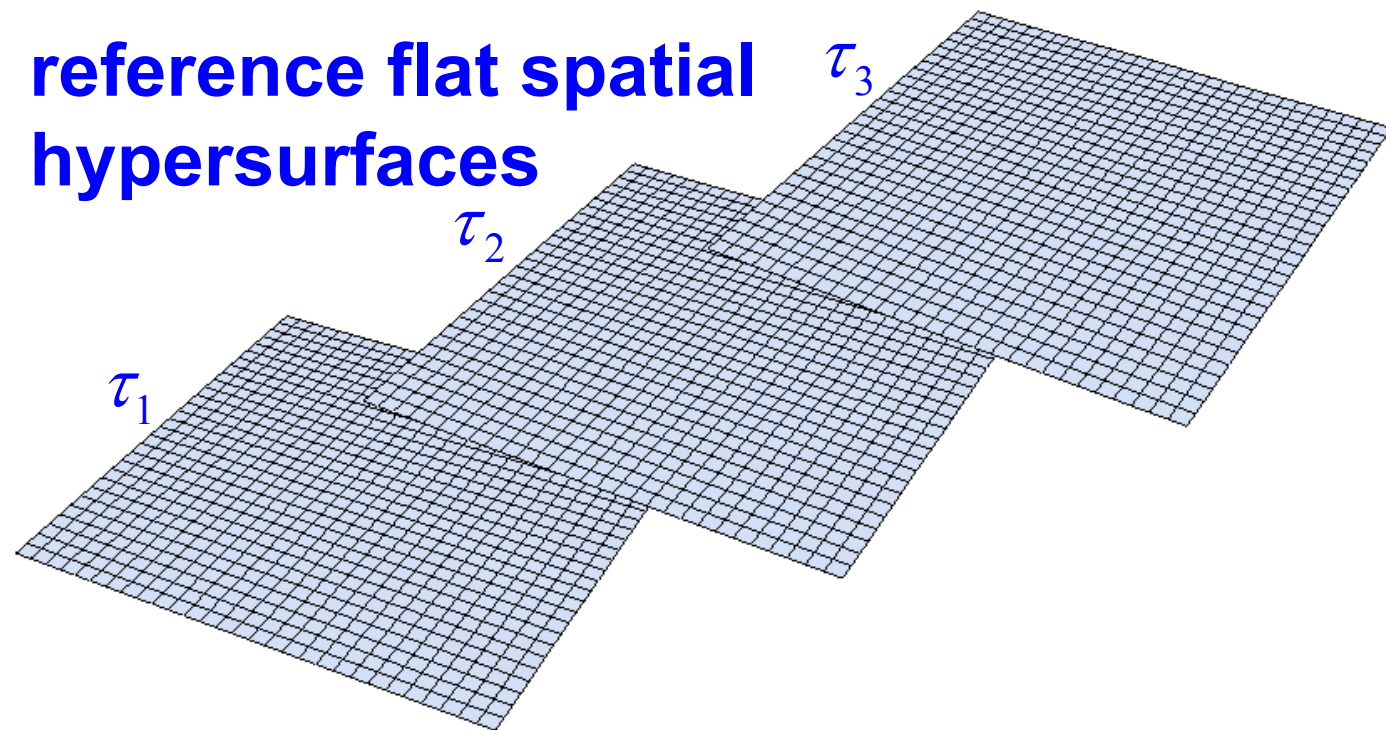
$$ds^2 = a^2(\tau) \left\{ d\tau^2 - \delta_{ij} dx^i dx^j \right\}$$

$\tau =$ conformal time

$$dt^2 = a^2(\tau) d\tau^2$$

Perturbed spacetime (10 degrees of freedom):

$$ds^2 = a^2(\tau) \left\{ (1 + \delta g_{00}) d\tau^2 - 2\delta g_{0i} d\tau dx^i - (\delta_{ij} + 2\delta g_{ij}) dx^i dx^j \right\}$$



scalar, vector, tensor decomposition

$$\delta g_{00} = 2A \qquad 1$$

$$\delta g_{0i} = S_i + \partial_i B \qquad 2+1$$
$$(\partial^i S_i = 0)$$

$$\delta g_{ij} = h_{ij} - \psi \delta_{ij} + \partial_i F_j + \partial_j F_i + \partial_i \partial_j E \qquad \frac{2+1+2+1}{\boxed{10}}$$
$$(h^i_i = 0 \ ; \ \partial^i h_{ij} = 0 \ ; \ \partial^i F_i = 0)$$

evolution of **scalar**, **vector**, and **tensor**
perturbations decoupled

Vector Perturbations:

- are not sourced by stress tensor
- decay rapidly in expansion

Tensor Perturbations:

- perturbations of transverse, traceless component of the metric: gravitational waves
- do not couple to stress tensor

Scalar Perturbations

- couple to stress tensor
- *density perturbations!*

Super-Hubble-radius

in synchronous gauge $A = B = 0$

and uniform Hubble flow gauge $B = E = 0$

$$\delta_+(t) = \delta_+(t_i) \left(t/t_i\right)^{2/3} \quad \text{matter-dominated}$$

$$\delta_+(t) = \delta_+(t_i) \left(t/t_i\right) \quad \text{radiation-dominated}$$

- in matter-dominated era

$$\delta_+(t) = \delta_+(t_i) \left(t/t_i\right)^{2/3} \quad \text{scales larger than Hubble radius}$$

$$\delta_+(t) = \delta_+(t_i) \left(t/t_i\right)^{2/3} \quad \text{scales smaller than Hubble radius}$$

- in radiation-dominated era

$$\delta_+(t) = \delta_+(t_i) \left(t/t_i\right) \quad \text{scales larger than Hubble radius}$$

$$\delta_+(t) = \text{constant} \quad \text{scales smaller than Hubble radius}$$

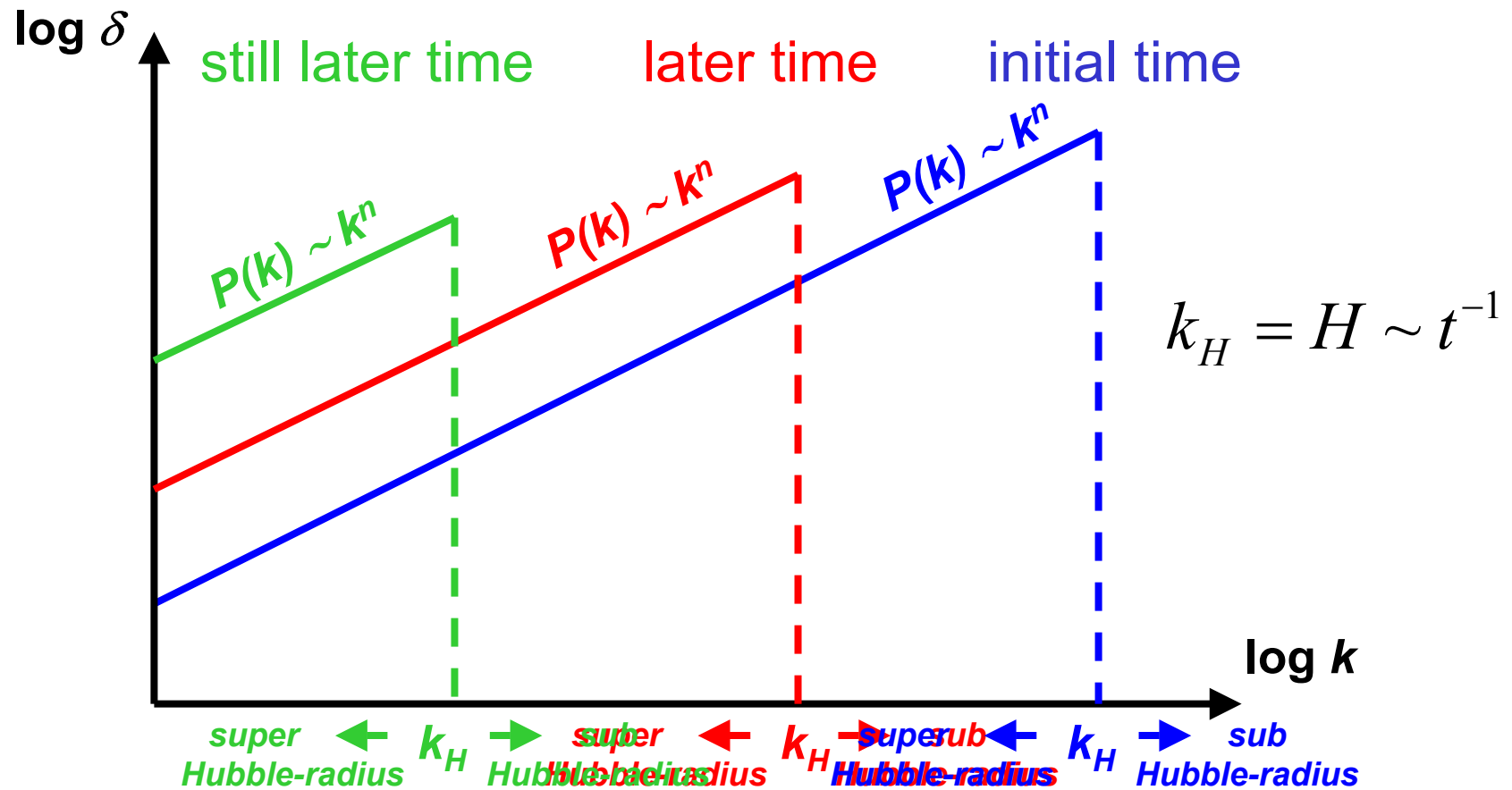
Harrison-Zel'dovich



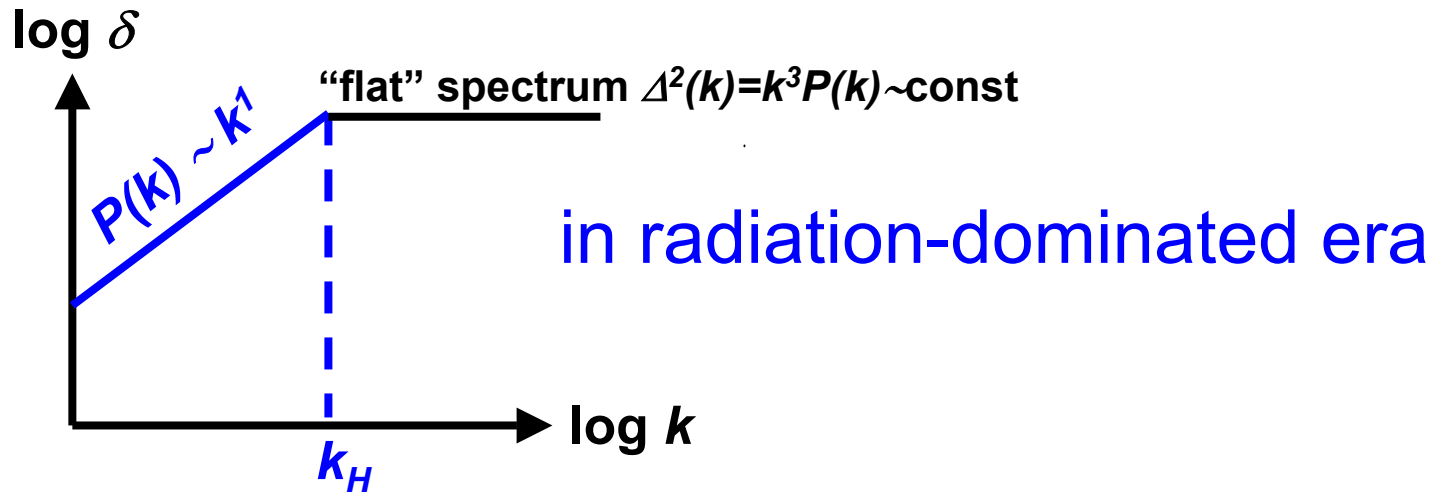
in radiation-dominated era

$\delta_+(t) = \delta_+(t_i) (t/t_i)$ scales larger than Hubble radius

$\delta_+(t) = \text{constant}$ scales smaller than Hubble radius

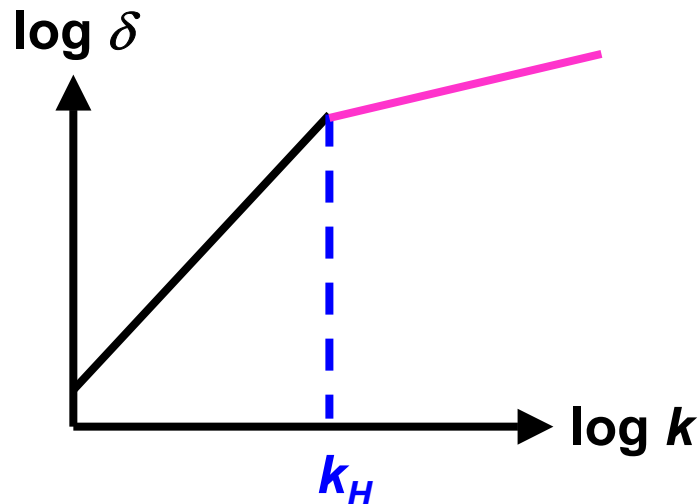


Harrison-Zel'dovich



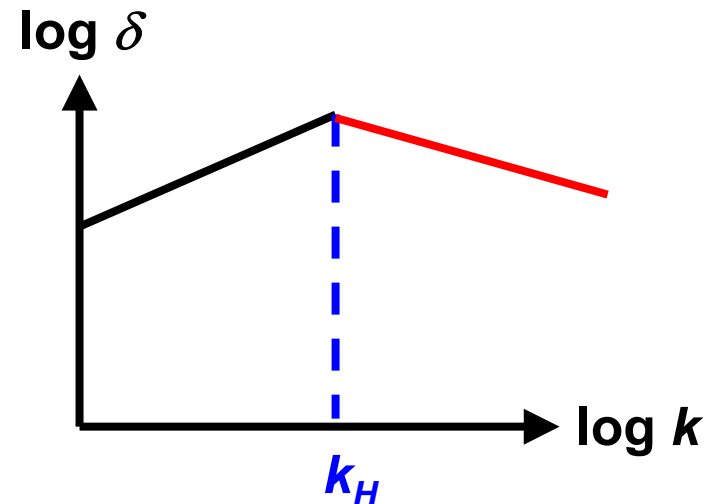
$$P(k) \propto k^n \quad n > 1$$

ultraviolet catastrophe

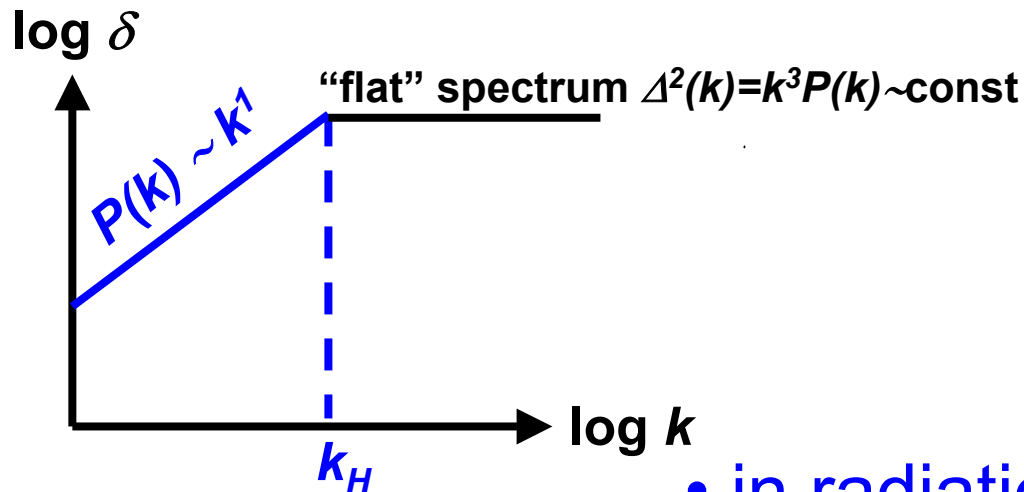


$$P(k) \propto k^n \quad n < 1$$

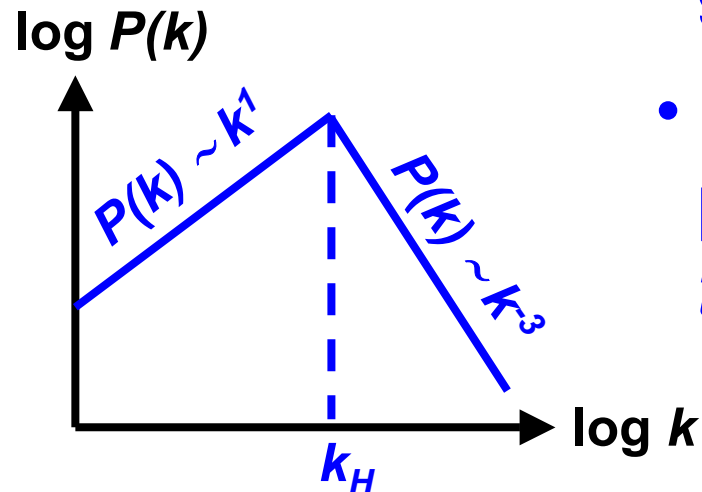
infrared catastrophe



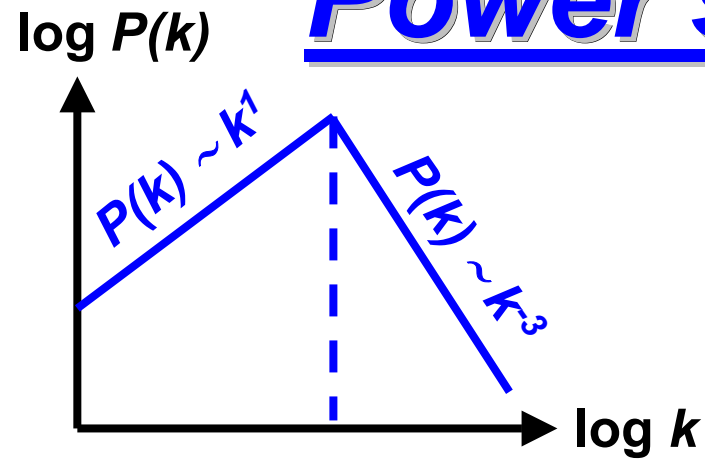
Harrison-Zel'dovich



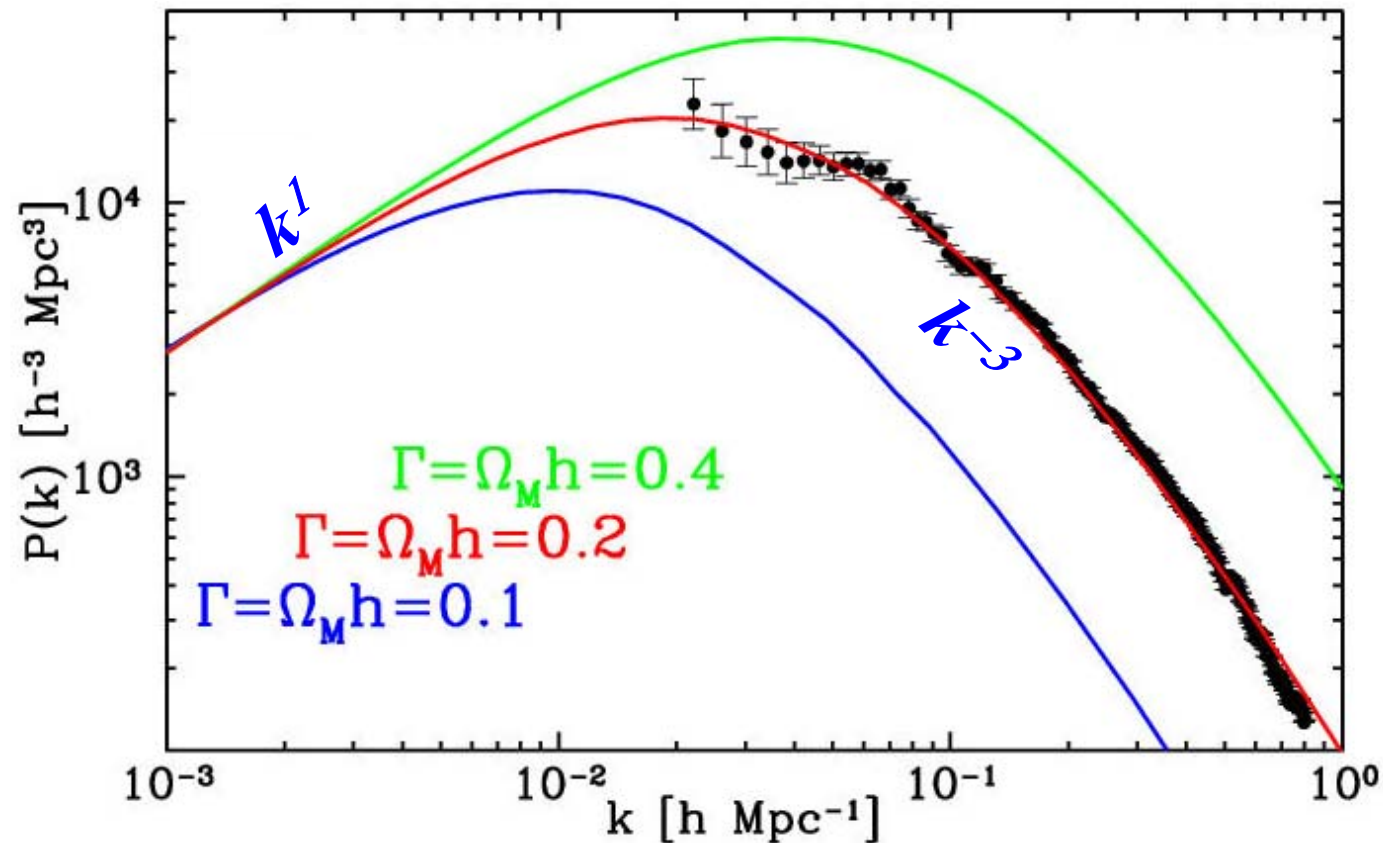
- in radiation-dominated era
no growth sub-Hubble radius
growth as t super-Hubble radius
- in matter-dominated era
power spectrum grows as $t^{2/3}$ on all scales



Power spectrum for CDM



matter-radiation equality



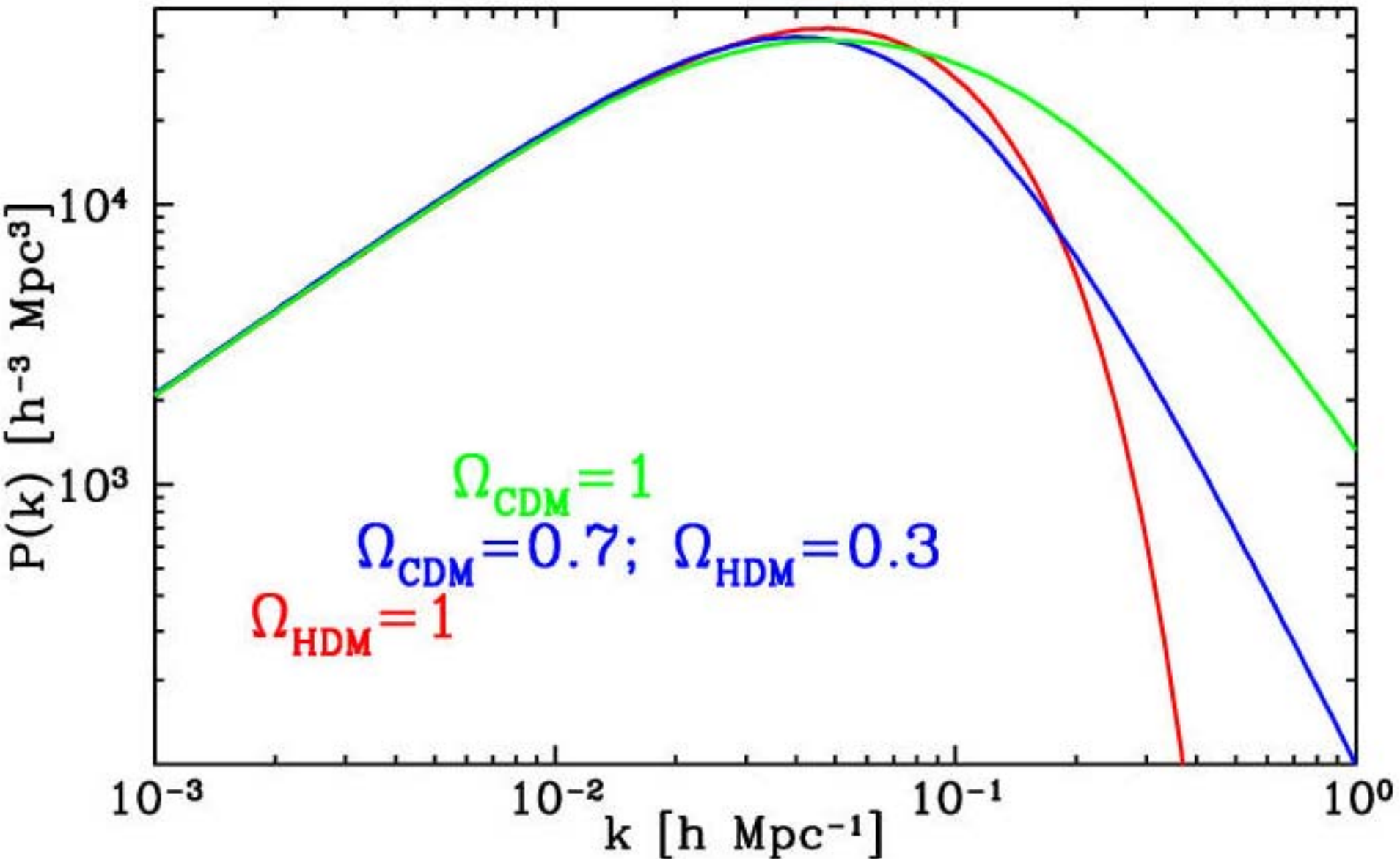
Dissipative processes

1. Collisionless phase mixing – free streaming

If dark matter is relativistic or semi-relativistic particles can stream out of overdense regions and smooth out inhomogeneities. The faster the particle the longer its free-streaming length.

Quintessential example: eV-range neutrinos

The evolved spectrum



Dissipative processes

1. Collisionless phase mixing – free streaming

If dark matter is relativistic or semi-relativistic particles can stream out of overdense regions and smooth out inhomogeneities. The faster the particle the longer its free-streaming length.

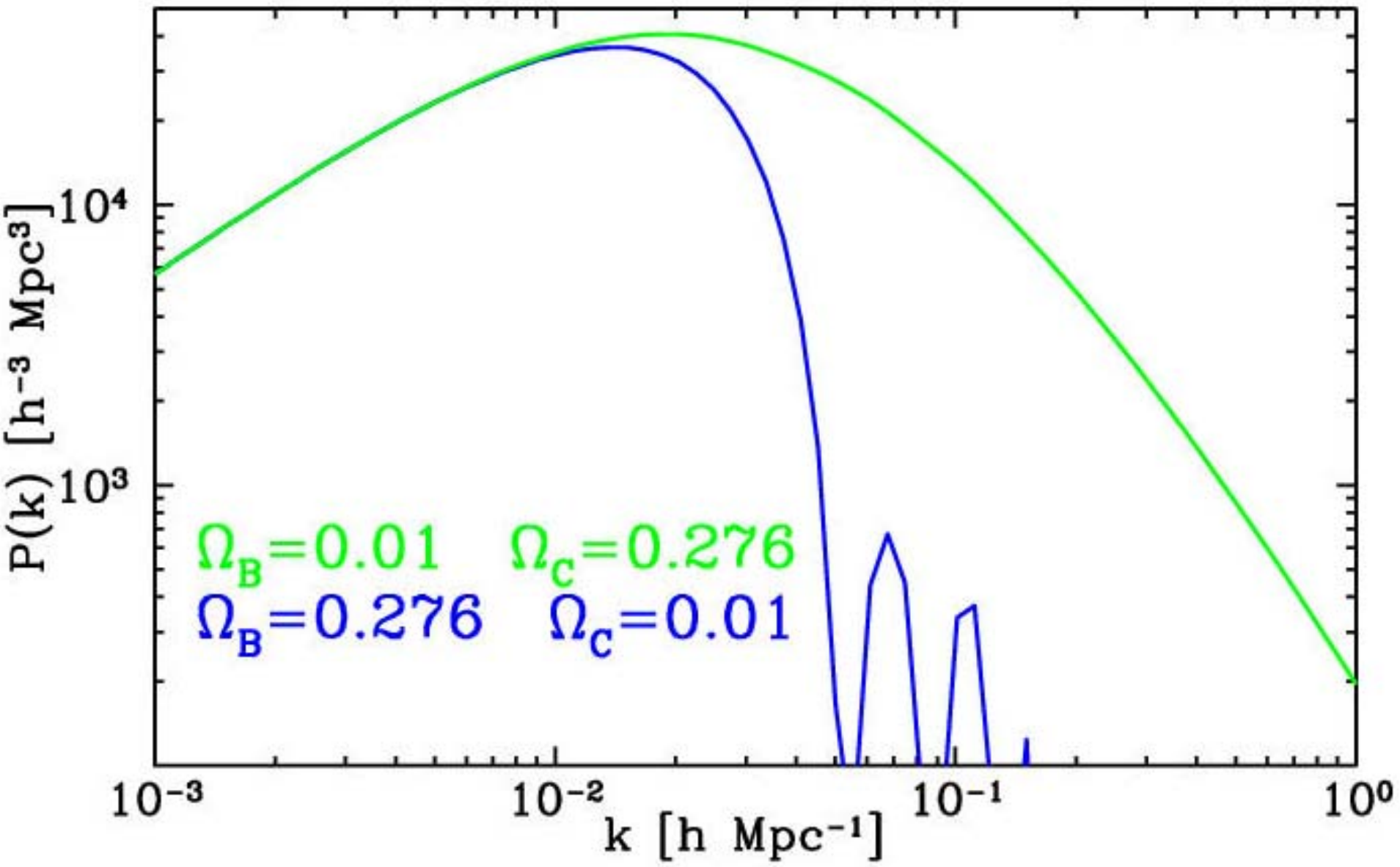
Quintessential example: eV-range neutrinos

2. Collisional damping – Silk damping

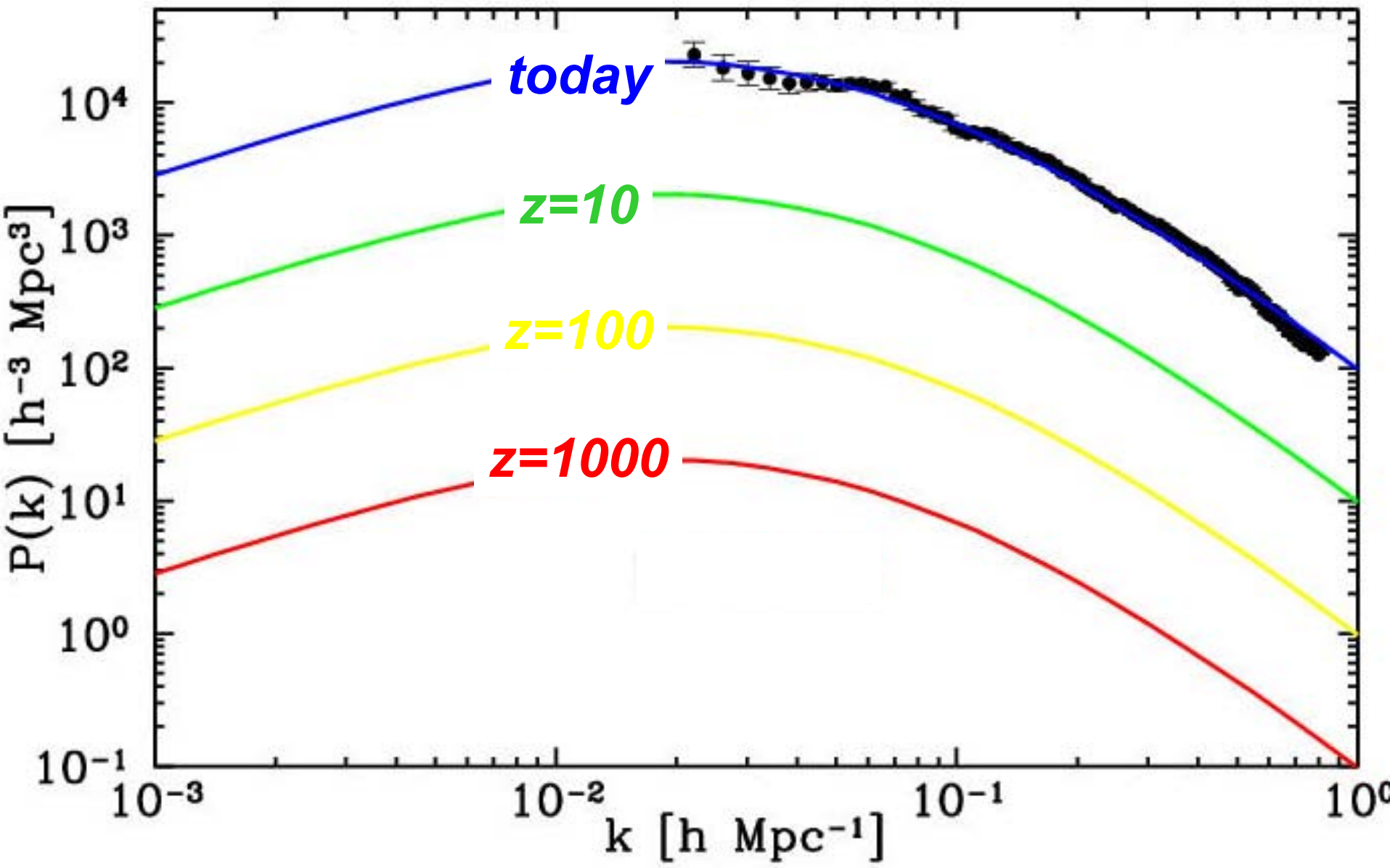
As baryons decouple from photons, the photon mean-free path becomes large. As photons escape from dense regions, they can drag baryons along, erasing baryon perturbations on small scales.

Baryon-photon fluid suffers damped oscillations.

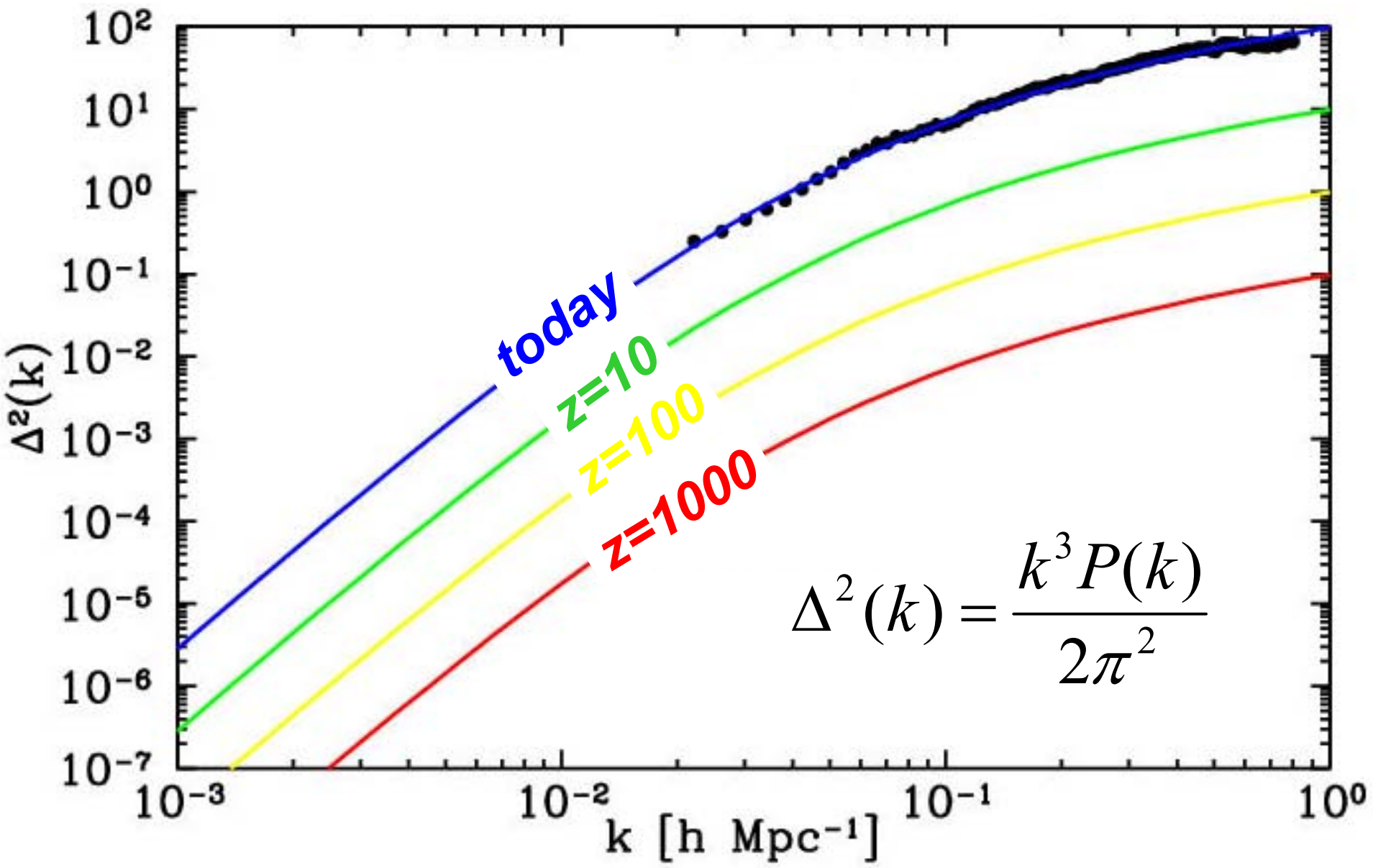
The evolved spectrum



Linear evolution

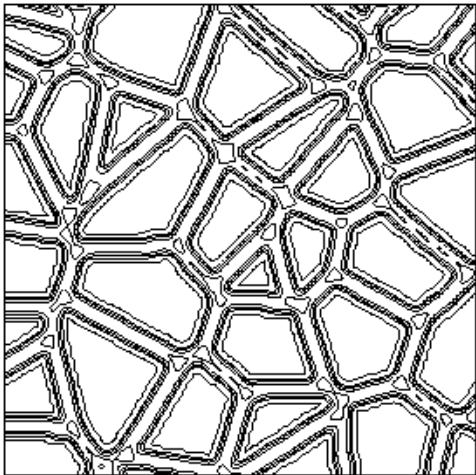


Linear evolution

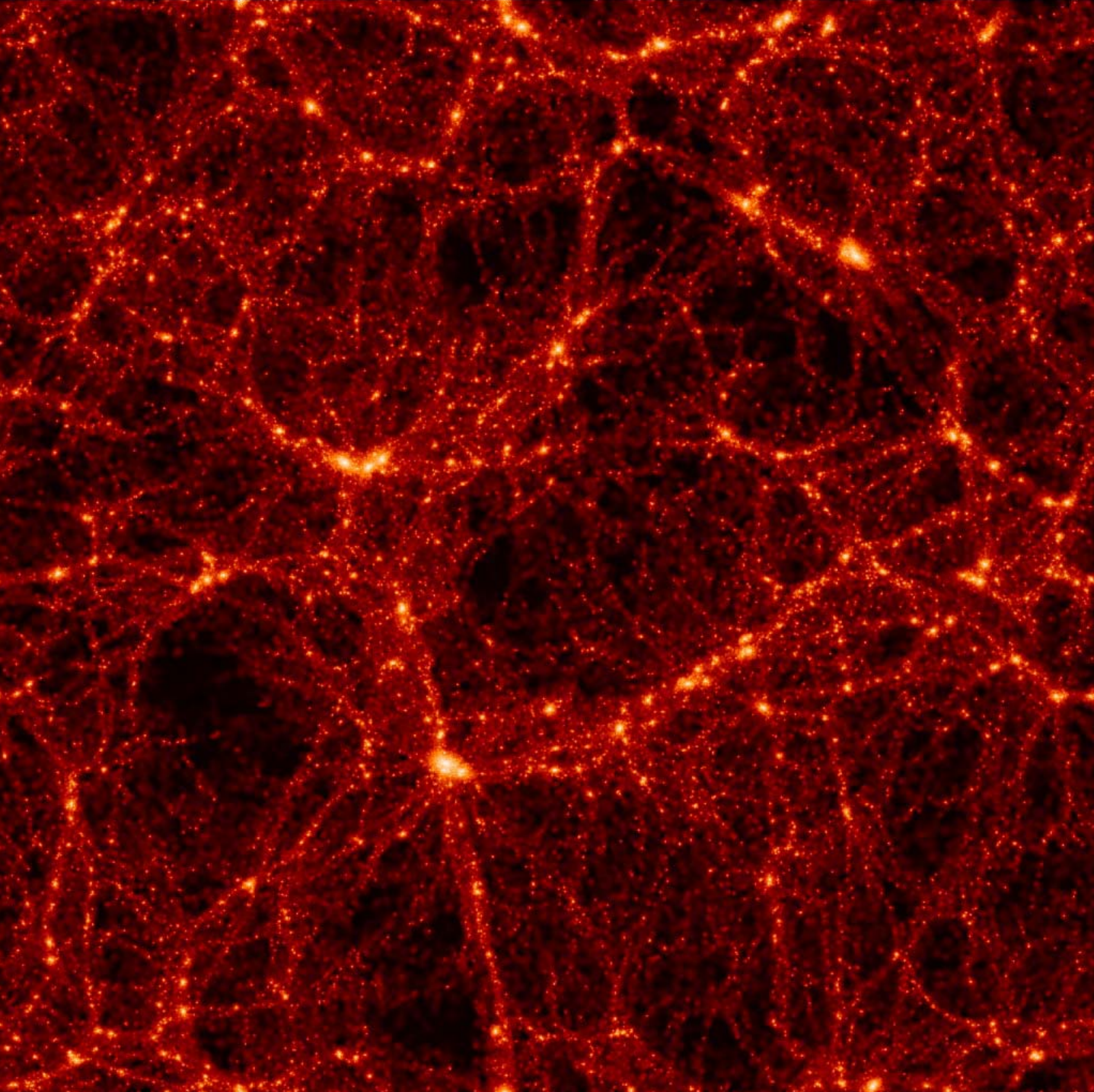


Life ain't linear!

- Many scales become nonlinear at about the same time
- Mergers from many smaller objects while larger scales form
- N-body simulations for dissipation-less dark matter
- Hydro needed for baryons
- Power spectrum well fit if $\Gamma = \Omega h \sim 0.2$
- There is more to life than the power spectrum

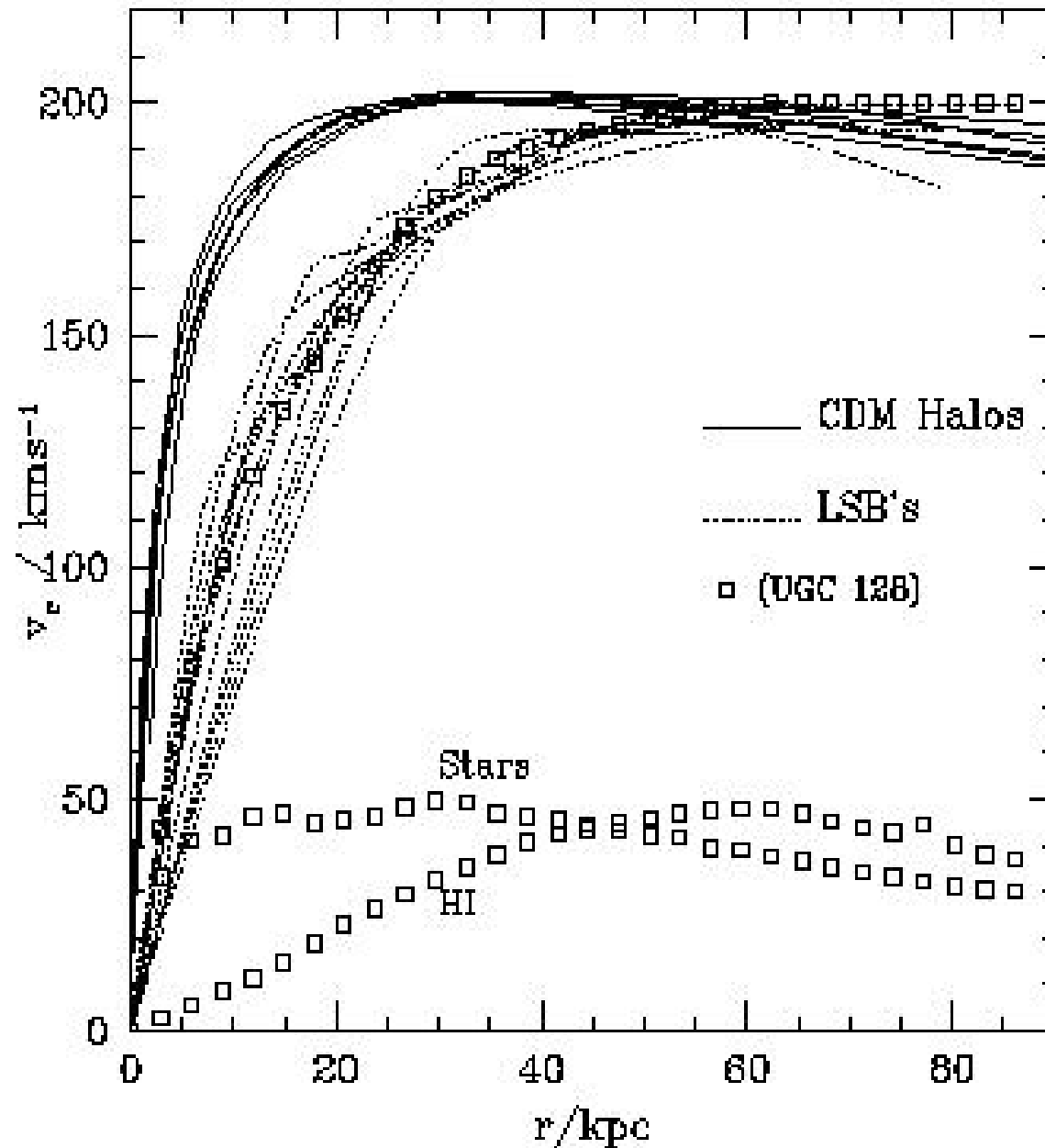


Alex Szalay



Large-
scale
structure
fits
well

Small-scale structure-cusps



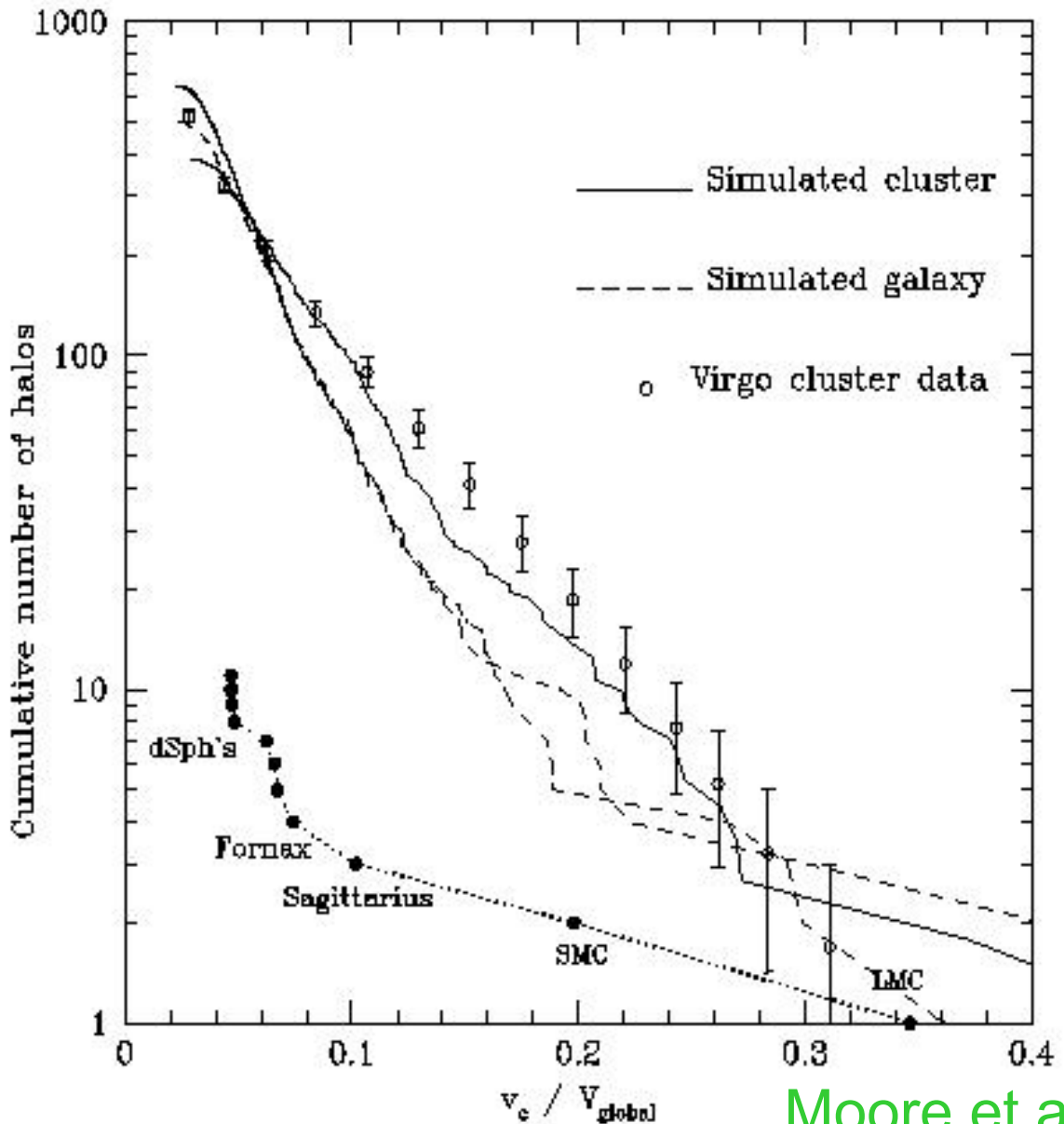
Small-scale structure-satellites

Cluster $5 \times 10^{14} M_{\odot}$

2 Mpc

Galaxy $2 \times 10^{12} M_{\odot}$

300 kpc



Rocky II: Growth of structure

- Linear regime: quantitative analysis

Jeans analysis

Sub-Hubble-radius perturbations (Newtonian)

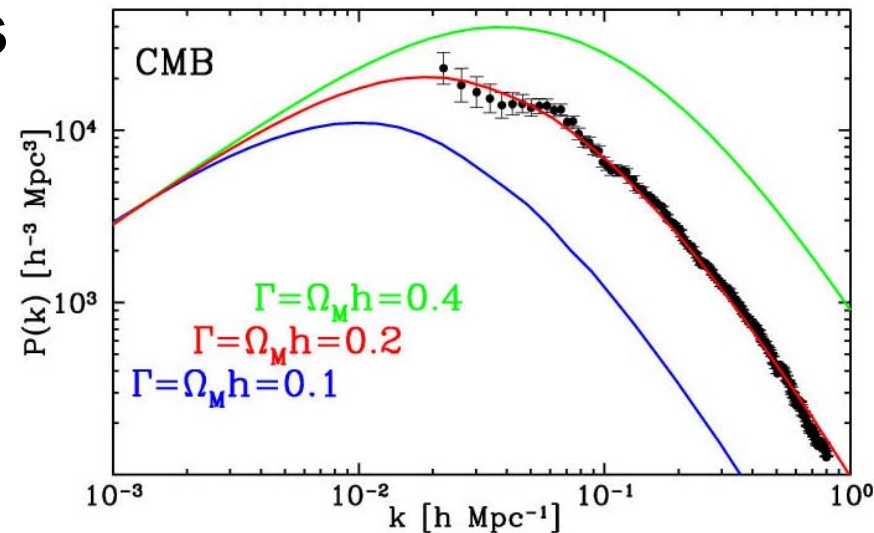
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